

A COMMENT ON THE RADIATIVE EFFICIENCY OF AGN

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ABSTRACT

A recent study of the accretion efficiency of the PG sample AGN, based on the thin accretion disk emission model, finds the accretion radiative efficiency is correlated with the black hole mass ($\eta \propto M_{\text{BH}}^{0.5}$). A followup study suggests the correlation is an artifact induced by selection effects. Here we point out there are two independent effects. The first is a sample selection effect, which leads to a high L/L_{Edd} AGN sample. The second effect is the observed small spread in the SED shape, which is not induced by selection effects. The second effect is what leads to the $\eta \propto M_{\text{BH}}^{0.5}$ relation in the PG sample. The physical reason for the small spread in the SED shape is an intriguing open question.

Thin accretion disk emission models allow to measure the absolute accretion rate in AGN, based on the optical luminosity L_{opt} , if the black hole mass M_{BH} is known. For example, Davis & Laor (2011, hereafter DL) derive

$$\dot{M} = 3.5 M_{\odot} \text{ yr}^{-1} (L_{\text{opt},45})^{3/2} M_8^{-0.89}, \quad (1)$$

where $M_8 = M_{\text{BH}}/10^8 M_{\odot}$, $L_{\text{opt}} \equiv \nu L_{\nu}$ at 4686 Å and $L_{\text{opt},45} = L_{\text{opt}}/10^{45} \text{ erg s}^{-1}$. The radiative efficiency, η , is then

$$\eta = \frac{L_{\text{bol}}}{\dot{M} c^2}. \quad (2)$$

where L_{bol} is the bolometric luminosity. DL applied this method to the complete and well defined PG quasar sample, using 80 of the 87 $z < 0.5$ AGN from the sample with UV observation. Single epoch broad line region based estimates of M_{BH} are available for all objects. The results shows that η tends to rise with M_{BH} . A best fit relation gives

$$\eta = 0.089 M_8^{0.52}. \quad (3)$$

The PG quasar sample is selected based on color, presence of broad permitted lines, and point-like morphology. The morphology criterion implies $L_{\text{opt}}/L_{\text{host}} > 1$. In addition, the following relations hold, $L_{\text{host}} \propto M_{\text{host}}$, $M_{\text{host}} \geq M_{\text{bulge}}$, where M_{host} and M_{bulge} are the host mass and bulge mass, and also $M_{\text{bulge}} \propto M_{\text{BH}}$. Thus, the morphology criterion implies a lower limit on $L_{\text{opt}}/M_{\text{BH}}$, and therefore a lower limit on $L_{\text{bol}}/L_{\text{Edd}}$, if $L_{\text{bol}} \propto L_{\text{opt}}$. Indeed, most PG quasars have $L_{\text{bol}}/L_{\text{Edd}} > 0.1$ (e.g. Fig.13 in DL). In contrast with some other AGN surveys, which can be significantly less contaminated by the host emission, like X-ray surveys, and therefore extend to much lower $L_{\text{bol}}/L_{\text{Edd}}$ values.

Raimundo et al. (2011, hereafter R11) used the thin accretion disk method to derive η for a small sample of AGN, part of which are nearby Seyfert galaxies. These objects have a distribution of M_{BH} values similar to the PG sample, but L_{opt} typically an order of magnitude lower (R11, Fig.17). The implied \dot{M} is a factor of ~ 30

lower than in the PG sample (R11, Fig.12), as expected from equation 1 above. The implied typical η values are ~ 3 times higher than for the PG quasars at similar M_{BH} (R11, Fig.16), as expected from equation 2 above, if $L_{\text{opt}}/L_{\text{bol}}$ has the same characteristic ratio as in the PG sample. R11 also analyzed a handful of higher z SDSS quasars, similar in properties to the PG quasars, and derive η values similar to the PGs for similar M_{BH} values (R11, Fig.16). R11 study carefully the range of parameters covered by the PG sample, assuming a fixed ratio for $L_{\text{opt}}/L_{\text{bol}}$, and conclude that the η vs. M_{BH} relation found for the PG quasar sample is an artifact of the sample selection criteria.

Below we explain why the η vs. M_{BH} relation in the PG sample is not an artifact. There are two independent relations among the observables that generate the correlation, and each needs to be carefully understood. The first comes from the PG sample selection criterion. As explained above, this leads to a lower limit on $L_{\text{opt}}/L_{\text{Edd}}$, and thus to a lower limit on \dot{M} for a given M_{BH} . But, this selection effect alone does not produce an η vs. M_{BH} relation. The second, and essential relation is the observed small spread in $L_{\text{opt}}/L_{\text{bol}}$ (DL, Fig.12). Applying both relations, $L_{\text{opt}} \propto M_{\text{BH}}$ and $L_{\text{opt}} \propto L_{\text{bol}}$ in eqs.1 & 2 above, leads inevitably to a rise of η with M_{BH} (with a power of 0.39, in the absence of scatter). The observed correlation between η and M_{BH} is only an artifact if selection effects generate the $L_{\text{opt}} \propto L_{\text{bol}}$ relation.

Therefore, it is crucial to understand that the $L_{\text{opt}} \propto L_{\text{bol}}$ relation is not a selection effect of the PG sample. The value of L_{bol} is mostly set by the far UV to soft X-ray part of the SED, which is independent of L_{opt} . The value of L_{opt} should come predominantly from larger disk radii than the UV and X-rays, and is not expected to provide a measure of L_{bol} . The expected SEDs of the PG sample, if all had a fixed η (~ 0.1), are shown in DL Fig.7. There is no selection effect which prevents the PG quasars from showing the predicted fixed η SED. Yet, at low M_{BH} the observed SED is systematically colder than the expected fixed η SED, while at the highest M_{BH} the observed SED is systematically hotter. The $\eta = 0.1$ SED is expected to peak at ~ 10 Rydberg for the $M_{\text{BH}} < 10^7 M_{\odot}$ PG sample objects, and at < 1 Rydberg for the $M_{\text{BH}} > 10^9 M_{\odot}$ objects. But, the SED generally peaks at 1 to a few Rydberg, regardless of M_{BH} .

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Low M_{BH} AGN should have hotter accretion disks than high M_{BH} AGN. But they don't. Why is this so? If η is set by the black hole spin a , then the rise of η with M_{BH} implies a rise of a with M_{BH} . In that case the innermost disk radius r_{in} , in units of $Rg = GM_{\text{BH}}/c^2$, shrinks with increasing M_{BH} , compensating for the drop in the disk temperature at a given radius (in Rg units). However, it is not clear why the two independent effects, the drop in r_{in} , and the rise in M_{BH} , happen to cancel

out and produce a relatively uniform SED (as far as one can tell).

Another option is that the accretion disk structure differs from the assumption of the thin disk model, due to some unmodelled physical processes, which produce a typical characteristic SED. In that case, η cannot be used to derive a . However, the inferred η relation may still provide some useful clue to processes in the innermost part of the accretion disk.

REFERENCES

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